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(54) **Single-transition-to-single-transition coherence transfer by cross-polarization: quantitation of cross-correlation effects in NMR**

(57) A method of NMR spectroscopy for achieving selective single-transition-to-single-transition cross-polarization (ST-CP) between two selected transitions in systems with two or more coupled spins comprises the following steps:

- 1) Choose RF pulses suitably shaped for a selective excitation experiment with an RF amplitude that is weaker than the scalar coupling constant J_{IS} of the I and S spin systems;
- 2) Simultaneously stimulate a pair of connected transitions of the I and S spin systems;
- 3) Record the response for the said selective excitation experiment and observe different pairs of connected transitions.

The method proposed is quite general provided the selectivity is sufficient to consider the relevant three-level subsystem to be isolated from all other energy levels. The method will be particularly useful to study differential relaxation of different multiplet components.

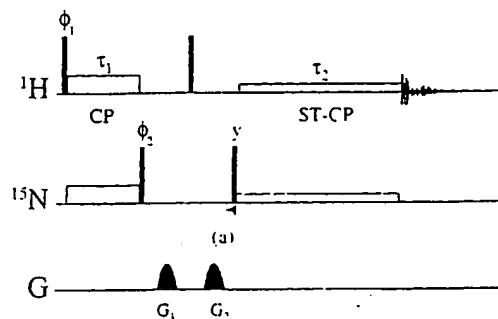


Figure 4

EP 1 098 204 A1

Description

[0001] The invention refers to a method of nuclear magnetic resonance (=NMR) spectroscopy for achieving selective single-transition-to-single-transition cross-polarization (ST-CP) between two selected transitions in systems with two or more coupled spins by irradiating radio frequency (=RF) pulses to a sample in a homogeneous magnetic field in z-direction.

I. INTRODUCTION

[0002] Cross-correlated relaxation phenomena have become important for the characterization of isotopically labelled macromolecules by solution-state NMR. The measurement of cross-correlated relaxation rates has enabled the direct measurement⁽¹⁻⁴⁾ of bond angles subtended by pairs of internuclear bond vectors in proteins without need to resort to empirical Karplus-type relationships. These angular constraints may also be used in conformational studies of small ligands bound to macromolecules using transferred cross-correlated relaxation,⁽⁵⁻⁷⁾ in analogy to transferred NOE. Measurement of ¹⁵N relaxation rates in both laboratory and rotating frames have yielded insight into exchange effects⁽⁸⁾ and dynamic behaviour⁽⁹⁾ of proteins in solution, the primary sources of relaxation in proteins being the chemical shift anisotropy (CSA) of the heteronuclei ¹³C or ¹⁵N and dipole-dipole (DD) coupling with directly bound protons.

[0003] Structural and dynamic studies become problematic with increasing molecular size due to increased transverse relaxation rates. A number of methods have been introduced to circumvent these effects. Transverse relaxation rates of amide protons may be reduced by exchanging all carbon-bound protons by deuterium. Alternatively one may use the interference between the CSA and DD interactions which leads to differential line narrowing⁽¹⁰⁻¹²⁾ of the individual multiplet components. In transverse relaxation optimized spectroscopy (TROSY)⁽¹³⁾ this effect is exploited by selecting the slowly-relaxing component in uncoupled heteronuclear correlation experiments. Improved versions⁽¹⁴⁻¹⁶⁾ using single-transition-to-single-transition coherence transfer yield a $\sqrt{2}$ sensitivity enhancement and have been employed⁽¹⁷⁾ to measure ¹³C relaxation rates in RNA and DNA bases. In general, coherence transfer is achieved via a series of INEPT-like transfers which limit the size of the proteins that may be studied. An avenue⁽¹⁸⁾ to study macromolecules of increasing molecular mass employs a combination of cross-correlated relaxation and evolution under scalar coupling to achieve polarization transfer.

[0004] Mehring, Wolff and Stoll⁽¹⁹⁾ have shown that it is possible to transfer coherence in a three-level system (such as occurs for deuterium with a non-vanishing quadrupole interaction) by simultaneous irradiation of two connected single-quantum transitions. This idea has been extended⁽²⁰⁾ to achieve coherence transfer between allowed and forbidden single- and multiple-quantum coherences.

[0005] It is an object of the invention to present a general method for studying differential relaxation of different multiplet components.

[0006] With the approach according to the present invention, one may also transfer coherence by selective cross-polarization between two selected transitions in a four-level system consisting of scalar coupled spins $I = 1/2$ and $S = 1/2$ in isotropic solution. This is achieved by the following steps:

1) Choose RF pulses suitably shaped for a selective excitation experiment with an RF amplitude that is weaker than the scalar coupling constant J_{IS} of the I and S spin systems;

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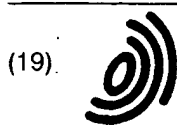
3) Record the response for the said selective excitation experiment and observe different pairs of connected transitions.

[0007] With the inventional method, selective single-transition-to-single-transition cross-polarization (ST-CP) can be achieved in heteronuclear scalar-coupled two-spin systems, such as amide ¹⁵N-¹H groups, by employing RF amplitudes that are much weaker than the scalar coupling constant J_{NH} , and by simultaneously irradiating a pair of connected transitions. Each of the four multiplet components in an uncoupled heteronuclear correlation spectrum may be observed in turn. The transfer benefits from differential line narrowing due to interference between the ¹⁵N chemical shift anisotropy and ¹⁵N-¹H dipolar coupling, which is clearly observed in ¹⁵N-labelled human ubiquitin at 600 MHz.

[0008] The spin system can be kept in isotropic or anisotropic solution during the experiment, depending on the desired environment of the sample material.

[0009] The intensities of the observed signals may depend upon the chosen pair of connected transitions because of cross-correlated relaxation effects.

[0010] An advantageous variant of the inventional method employs semi-selective cross-polarization during a preparation period τ_1 to transfer coherence through the scalar coupling from both doublet components of the selected



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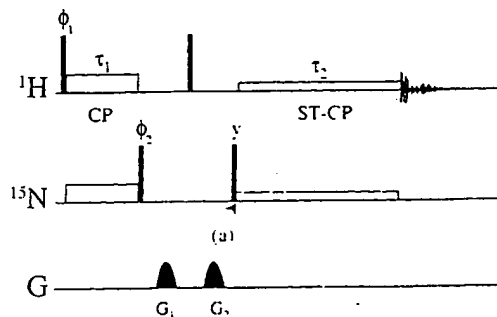


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[0006] With the approach according to the present invention, one may also transfer coherence by selective cross-polarization between two selected transitions in a four-level system consisting of scalar coupled spins $I = 1/2$ and $S = 1/2$ in isotropic solution. This is achieved by the following steps:

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[0010] An advantageous variant of the inventional method employs semi-selective cross-polarization during a preparation period τ_1 to transfer coherence through the scalar coupling from both doublet components of the selected

source spin I to the two doublet components of the directly bound second spin S.

[0011] This may be improved by using matched RF fields with amplitudes of the order of the heteronuclear coupling constant J_{IS} .

[0012] In a preferred variant, the method according to the present invention may be characterised by applying the following pulse sequence to the spin system I and S :

- a) irradiating a $\pi/2$ RF pulse acting on spins I with a phase-cycle $\Phi_1 = y, y, -y, -y$;
- b) applying two RF fields with matched amplitudes of $\omega_1/2\pi = \omega_1^S/2\pi$ and carrier frequencies positioned at the midpoints of the I and the S doublets during a preparation time period $\tau_1 = J_{IS}^{-1}$;
- c) irradiating a $\pi/2$ RF pulse acting on spins S with a phase-cycle $\Phi_2 = y, -y, -y, y$;
- d) destroying any residual transverse magnetization by applying a first spoil gradient pulse G_1 followed by a proton $\pi/2$ pulse and subsequently by applying a second spoil gradient pulse G_2 ;
- e) irradiating a $\pi/2$ pulse acting on the spins S;
- f) applying two RF fields with matched amplitudes of $\omega_1/2\pi = \omega_1^S/2\pi$ for single-transition-to-single-transition cross-polarization, on-resonance with any chosen pair of connected transitions during a second time period τ_2 , the amplitudes of the RF fields being reduced with respect to the amplitudes in step b)
- g) using a receiver phase-cycling $\Phi_{rec} = x, -x, x, -x$.

[0013] This can be further improved by inserting a suitable relaxation delay at a time point (a) before the ST-CP interval τ_2 to monitor the relaxation behaviour of the spin S coherences whilst observing the spin I doublet components.

[0014] Another improvement of this method is characterised by inserting a Carr-Purcell-Meiboom-Gill (CPMG) multiple refocussing sequence at time point (a) before the ST-CP interval τ_2 to monitor the relaxation behaviour of the spin S coherences whilst observing the spin I doublet components.

[0015] It may be also of advantage using a Carr-Purcell-Meiboom-Gill (CPMG) spin echo sequence for T_2 measurements whereby the intervals between the π pulses of the CPMG pulse train are set to $2\tau = n/J_{IS}$, where n is an integer.

[0016] Alternatively, the method according to the present invention may be modified by inserting a spin-lock field at time point (a) before the ST-CP interval τ_2 for $T_{1\rho}$ measurements.

II. THEORY

[0017] Figure 1 shows the energy-level diagram for a scalar coupled two-spin system IS. In analogy to the case of deuterium,^(19,20) it is possible to transfer coherence by simultaneously irradiating two connected transitions. Spin-state selective coherence transfer in a four-level system relies on irradiating the two transitions in a highly selective manner. In an ideal case, the presence of the fourth energy level should not affect the cross-polarization pathway of interest.

[0018] In the theory outlined below we consider cross-polarization from the coherence $I_x^{(1,2)}$ associated with a nitrogen transition to the progressively connected proton coherence $I_x^{(2,4)}$. In the presence of two RF fields which are on-resonance for the selected proton and nitrogen transitions the Hamiltonian in the laboratory frame may be written as

$$\begin{aligned}
 H^{lab} = & \omega_0^I I_z + \omega_0^S S_z + 2\pi J_{IS} I_z S_z \\
 & + \omega_1^I \cos(\omega_0^{(2,4)} t) I_x + \omega_1^I \sin(\omega_0^{(2,4)} t) I_y \\
 & + \omega_1^S \cos(\omega_0^{(1,2)} t) S_x + \omega_1^S \sin(\omega_0^{(1,2)} t) S_y, \quad (1)
 \end{aligned}$$

where the RF amplitudes are given by $\omega_1^I = -\gamma^I B_1^I$ and $\omega_1^S = -\gamma^S B_1^S$. For clarity the initial RF phases have been assumed to be zero. It is convenient to recast this Hamiltonian in terms of single transition operators.⁽²¹⁻²²⁾ The proton and nitrogen doublets, respectively, are given by

$$I_x = I_x^{(1,3)} + I_x^{(2,4)} \text{ and } S_x = I_x^{(1,2)} + I_x^{(3,4)}. \quad (2)$$

[0019] The resonance frequencies of these transitions result from the Zeeman interaction and the scalar coupling. We define an interaction representation using the transformation⁽²³⁾

$$U_0 = \exp(-iH_0 t)$$

where $H_0 = \omega_0^I I_z + \omega_0^S S_z + 2\pi J_{IS} I_z S_z$. The laboratory frame Hamiltonian of Eq. (1) becomes

$$\begin{aligned} H^{ref} = & \omega_1^I I_x^{(2,4)} + \omega_1^S I_x^{(1,2)} \\ & + \omega_1^I I_x^{(1,3)} \cos(2\pi J_{IS} t) + \omega_1^I I_y^{(1,3)} \sin(2\pi J_{IS} t) \\ & + \omega_1^S I_x^{(3,4)} \cos(2\pi J_{IS} t) + \omega_1^S I_y^{(3,4)} \sin(2\pi J_{IS} t) \end{aligned} \quad (3)$$

[0020] For very weak RF fields ($\omega_1^I, \omega_1^S \ll 2\pi J_{IS}$) one may neglect the oscillating components by invoking the secular approximation. Thus, as an initial approximation, we consider only the first two terms in the Hamiltonian. These terms are responsible for on-resonance coherence transfer. We shall see later that the other terms can lead to off-resonance coherence transfer if the amplitude of the spin-lock fields is not much weaker than the coupling constant.

[0021] Figure 2 shows the evolution of the density operator under the influence of the Hamiltonian in Eq. (3), truncated to include only the first two terms. Cross-polarization is described by a rotation of the initial vector $\sigma(0) = I_x^{(1,2)}$ onto $I_x^{(2,4)}$ about an effective field

$$\omega_{eff} = \sqrt{(\omega_1^I/2)^2 + (\omega_1^S/2)^2}.$$

This amounts to a rotation confined to a three-dimensional subspace of the full sixteen-dimensional Liouville space,

$$\{2I_x^{(1,2)}, 2I_x^{(2,4)}, 2I_y^{(1,4)}\} \quad (4)$$

[0022] The evolution of the density operator in the case of matched RF fields is given by.

$$\sigma(t) = \frac{1}{2} (1 + \cos \omega_{eff} t) I_x^{(1,2)} + \frac{1}{2} (1 - \cos \omega_{eff} t) I_x^{(2,4)} - \frac{1}{\sqrt{2}} (\sin \omega_{eff} t) I_y^{(1,4)} \quad (5)$$

[0023] In contrast to conventional semi-selective cross-polarization⁽²⁵⁾, the transfer does not depend on the coupling constant J_{IS} . The efficiency of the transfer is a maximum when the Hartmann-Hahn condition is fulfilled ($\omega_1^I = \omega_1^S$). It should be noted that

$$\text{tr}(I_x^{(2,4)} I_x^{(2,4)}) = 1/2,$$

corresponding to one component of the doublet. In the absence of relaxation, the maximum expectation value of the observable

$$\langle I_x^{(2,4)} \rangle$$

is therefore 1/2, which occurs for a spin-lock duration of $\tau = \pi/\omega_{eff}$. Thus RF field amplitudes $\omega_1^I/2\pi = \omega_1^S/2\pi = 26$ or 11

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is therefore 1/2, which occurs for a spin-lock duration of $\tau = \pi/\omega_{eff}$. Thus RF field amplitudes $\omega_1^I/2\pi = \omega_1^S/2\pi = 26$ or 11

Hz yield optimum transfer times of 27 or 64 ms respectively.

[0024] One should note that coherence transfer processes $(1,2) \rightarrow (2,4)$ or $(1,3) \rightarrow (3,4)$ between progressively connected transitions in the two subsystems $(1,2,4)$ and $(1,3,4)$ occur via the double quantum transition $f_y^{(1,4)}$ whereas coherence transfer processes $(1,2) \rightarrow (1,3)$ or $(3,4) \rightarrow (2,4)$ between regressively connected transitions in the two subsystems $(1,2,3)$ and $(2,3,4)$ occur via the zero quantum transition $f_y^{(2,3)}$.

[0025] Equation (5) gives us insight into the overall time dependence of the density operator. However, in systems such as ^{15}N - H^{N} in proteins, the scalar coupling is usually of the order of $J_{\text{NH}} = 95$ Hz, so that the secular approximation may be violated and the oscillating components in the Hamiltonian may have to be considered. Computer simulations have been carried out using subroutines of GAMMA⁽²⁴⁾ and the complete Hamiltonian of Eq. (3). Although the dominant feature is due to the transfer function given in Eq. (5), the signals are modulated by additional frequencies. A frequency component corresponding to the coupling constant $2\pi J_{\text{IS}}$ represents a parasitic transfer (light dashed arrows in Fig. 1) which is of the order of 5 % of the maximum intensity of the on-resonance transfer. A second, more pronounced, low-frequency modulation may also be observed. The frequency of this modulation increases with increasing RF field amplitude. It is found that this additional modulation does not have a detrimental effect on the efficiency of cross-polarization provided $\omega_1/2\pi \leq J_{\text{IS}}/3$.

III. RELAXATION

[0026] Single-transition-to-single-transition cross-polarization provides interesting insight into relaxation occurring during transfer. We have chosen a phenomenological approach by constructing a Liouvillian superoperator which includes coherent interactions given by the Hamiltonian in Eq. (3) and relaxation. The relaxation superoperator is diagonal if cross-relaxation is neglected, with diagonal elements composed of the self-relaxation rates of each coherence.

[0027] Figure 3 shows the trajectory of the density operator as a function of time. The relaxation times are representative of an amino acid in the rigid core of a small protein. All density operator components are attenuated by relaxation, and the vector in Fig. 3 follows a helical trajectory around ω_{eff} . One may note that, even for a protein, the efficiency of cross-polarization for a contact time $\tau = 24$ ms with RF field amplitudes of 26 Hz is about 68 % of the amplitude expected in the absence of relaxation.

[0028] It is also important to note that, due to the spin-state selective nature of the transfer, the efficiency will be governed by the individual relaxation rates for each coherence. Thus cross-polarization between slowly relaxing components will be more efficient than between rapidly relaxing components. This is another manifestation of the so-called TROSY effect.

IV. EXPERIMENT

[0029] The preparation period τ_1 in Fig. 4 employs semi-selective cross-polarization⁽²⁵⁻²⁶⁾ to transfer coherence through the scalar coupling from both doublet components of a selected source spin H^{N} to the two doublet components of the directly bound ^{15}N . This may be achieved using matched fields with amplitudes of the order of the heteronuclear coupling to overcome problems with RF field inhomogeneity⁽²⁵⁾. We subsequently rotate the resulting S_x magnetization along the z axis and destroy any residual transverse magnetization, notably that of water, with two spoil gradients and a proton pulse. After another $\pi/2$ applied to ^{15}N we have both nitrogen doublet components $S_x = f_x^{(1,2)} + f_x^{(3,4)}$ at time point (a). We then apply two low-amplitude fields for single-transition-to-single-transition cross-polarization, on-resonance with any chosen pair of connected transitions. This yields a spectrum composed of a proton doublet whereby the on-resonance doublet component displays the build-up behaviour given by Eq. (5). The other off-resonance doublet component is modulated as a function of τ_2 at a frequency corresponding to the coupling constant. For example, if one chooses the on-resonance transfer $f_x^{(1,2)} \rightarrow f_x^{(2,4)}$ then the off-resonance transfer would be $f_x^{(3,4)} \rightarrow f_x^{(1,3)}$, corresponding to the component which is diagonally opposed in the HSQC spectrum.

[0030] One can modify the pulse sequence in Fig. 4 to measure $T_{1\rho}$ or T_2 relaxation times for the individual nitrogen doublet components. If one inserts a suitable relaxation delay at time point (a) in Fig. 4 then the relaxation behaviour of the nitrogen coherences may be monitored whilst observing the proton doublet components. For $T_{1\rho}$ measurements a spin-lock field can be inserted. For T_2 measurements a Carr-Purcell-Meiboom-Gill (CPMG) spin echo sequence can be used. The intervals between the π pulses of the CPMG pulse train are set to $2\tau = n/J_{\text{IS}}$ where n is an integer. This averages the auto-relaxation rates of in-phase S_x and antiphase $2S_x I_z$ coherences as the scalar coupling remains active during these intervals.⁽²⁷⁾ In this case the effective relaxation matrix reduces to diagonal form with each doublet component relaxing mono-exponentially at a rate $R_{av} \pm \delta$ where

$$R_{av} = (\rho_N + \rho_{\text{NH}})/2,$$

δ is the CSA-dipole cross correlated relaxation rate and ρ_N and ρ_{NH} are the autorelaxation rates of in- and antiphase terms S_x and $2S_xI_z$. In the case of T_1 measurements, where there is no evolution under the scalar coupling, the relaxation will in general become multiexponential.⁽²⁸⁾

[0031] A 100 mM sample of ^{15}N enriched boc-protected glycine in DMSO was used to illustrate the transfer process.

This is a convenient test sample and may be considered as an isolated ^{15}N - ^1H fragment ($J_{NH} \approx 94$ Hz) within a small molecule with correspondingly long T_2 's. All experiments were performed on a Bruker DRX 600 MHz spectrometer with a TBI probe equipped with triple axis gradients. Figure 5 shows the signal intensity due to on-resonance ST-CP as a function of the spin lock duration τ_2 . For clarity only two curves are displayed which correspond to coherence transfer from the high-frequency nitrogen doublet component $I_x^{(1,2)}$ to one of the proton doublet components. Since the T_2 's of the two proton doublet components are almost equal, the differences in the two curves are primarily due to the fact that the density operator passes through states with different relaxation times, the progressive connectivity being mediated via double quantum coherence whereas the regressive connectivity passes through zero quantum coherence. The other two cross-polarization paths from the low-frequency nitrogen coherence $I_x^{(3,4)}$ (not shown) are qualitatively similar. Also displayed is a simulation of the evolution of the density operator, using a Liouvillian superoperator composed of the sum of coherent and relaxation superoperators. Relaxation times $T_{1\rho}^S$ and $T_{1\rho}^I$ for the nitrogen and proton transitions may be measured by modifying the pulse sequence, as described previously. The decay yields a value of $T_{1\rho}^S = 2.86$ s for both nitrogen transitions. If a proton spin-lock field is inserted immediately before acquisition, the decay yields a value of $T_{1\rho}^I = 383$ ms for both proton transitions. One may then vary the relaxation time of the zero- or double-quantum transition empirically until an agreement with the experiments is attained.

[0032] It is instructive to take a Fourier transform to unravel the different frequency components in the build-up curves of Fig. 5. This was carried out using MATLAB and is shown in Fig. 6. The most intense peak is centered at $\omega_2 = 18$ Hz corresponding to RF field strengths $\omega_1^S/2\pi = \omega_1^I/2\pi = 26$ Hz ($\omega_{eff}/2\pi = 18.4$ Hz). Additional fine structure may be observed giving evidence of a low-frequency modulation which arises due to a breakdown in the secular approximation. One may also observe a doublet centered at the coupling constant $J_{NH} = 94$ Hz corresponding to a residual 'parasitic' transfer from the other nitrogen doublet component. These effects are consistent with simulations. Also displayed in Fig. 6 are simulations in which the relaxation time T_2^{ZQC} of the zero quantum coherence $I_x^{(2,3)}$ is varied. The best fit may be observed for $T_2^{ZQC} \approx 30$ ms.

V. APPLICATION

[0033] The principles of ST-CP were applied to relaxation measurements in a 1.5 mM sample of ^{15}N -labeled human ubiquitin (VLI Research) in $\text{H}_2\text{O} : \text{D}_2\text{O} = 9 : 1$ buffered at pH = 4.5 with 20 mM perdeuterated acetic acid at a temperature of 303 K and a static field of 14 T.

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VI. CONCLUSIONS

[0035] It has been shown that one may achieve single-transition-to-single-transition coherence transfer in a scalar-coupled two-spin system by cross polarization. The method proposed is quite general provided the selectivity is sufficient to consider the relevant three-level subsystem to be isolated from all other energy levels. The method will be particularly useful to study differential relaxation of different multiplet components. The principles outlined here may cast new light on a variety of cross-correlated relaxation phenomena.

Figure Captions

[0036] FIG. 1. Energy-level diagram for a scalar-coupled two-spin system IS, such as is found in backbone amide N-H groups of ^{15}N labelled proteins and peptides. The nitrogen doublet components correspond to transitions (1,2) and (3,4). The proton doublet components correspond to transitions (1,3) and (2,4). Also shown is a schematic representation of single-transition-to-single-transition cross-polarization (ST-CP) transfer. The main pathways correspond to on-resonance (thick solid arrow) and off-resonance transfer (thick dashed arrow). Also shown are undesired parasitic transfer pathways (thin dashed arrows).

[0037] FIG. 2. The evolution of the density operator during single-transition-to-single-transition cross-polarization corresponds to a rotation within a three-dimensional subspace of the complete 16 dimensional Liouville space. Coher-

Hz yield optimum transfer times of 27 or 64 ms respectively.

[0024] One should note that coherence transfer processes $(1,2) \rightarrow (2,4)$ or $(1,3) \rightarrow (3,4)$ between progressively connected transitions in the two subsystems $(1,2,4)$ and $(1,3,4)$ occur via the double quantum transition $f_y^{(1,4)}$ whereas coherence transfer processes $(1,2) \rightarrow (1,3)$ or $(3,4) \rightarrow (2,4)$ between regressively connected transitions in the two subsystems $(1,2,3)$ and $(2,3,4)$ occur via the zero quantum transition $f_y^{(2,3)}$.

[0025] Equation (5) gives us insight into the overall time dependence of the density operator. However, in systems such as ^{15}N - H^{N} in proteins, the scalar coupling is usually of the order of $J_{\text{NH}} = 95$ Hz, so that the secular approximation may be violated and the oscillating components in the Hamiltonian may have to be considered. Computer simulations have been carried out using subroutines of GAMMA⁽²⁴⁾ and the complete Hamiltonian of Eq. (3). Although the dominant feature is due to the transfer function given in Eq. (5), the signals are modulated by additional frequencies. A frequency component corresponding to the coupling constant $2\pi J_{\text{IS}}$ represents a parasitic transfer (light dashed arrows in Fig. 1) which is of the order of 5 % of the maximum intensity of the on-resonance transfer. A second, more pronounced, low-frequency modulation may also be observed. The frequency of this modulation increases with increasing RF field amplitude. It is found that this additional modulation does not have a detrimental effect on the efficiency of cross-polarization provided $\omega_1/2\pi \leq J_{\text{IS}}/3$.

III. RELAXATION

[0026] Single-transition-to-single-transition cross-polarization provides interesting insight into relaxation occurring during transfer. We have chosen a phenomenological approach by constructing a Liouvillian superoperator which includes coherent interactions given by the Hamiltonian in Eq. (3) and relaxation. The relaxation superoperator is diagonal if cross-relaxation is neglected, with diagonal elements composed of the self-relaxation rates of each coherence.

[0027] Figure 3 shows the trajectory of the density operator as a function of time. The relaxation times are representative of an amino acid in the rigid core of a small protein. All density operator components are attenuated by relaxation, and the vector in Fig. 3 follows a helical trajectory around ω_{eff} . One may note that, even for a protein, the efficiency of cross-polarization for a contact time $\tau = 24$ ms with RF field amplitudes of 26 Hz is about 68 % of the amplitude expected in the absence of relaxation.

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[0039] FIG. 4. Pulse sequence used for coherence transfer by cross-polarization. The first semi-selective cross-polarization (CP) period has a duration $\tau_1 = J_{\text{IS}}^{-1} = 10.9$ ms and matched amplitudes of $\omega_1^I/2\pi = \omega_1^S/2\pi \approx 70$ Hz. The carrier frequencies are positioned at the midpoints of the proton and nitrogen doublets. During the second, selective single-transition-to-single-transition cross-polarization (ST-CP) period τ_2 , the matched amplitudes are reduced to $\omega_1^I/2\pi = \omega_1^S/2\pi \approx 26$ Hz and the carrier frequencies moved to any two connected transitions which appear in the undecoupled multiplet, e.g. $\omega_0^{(1,2)} = \omega_0^S + \pi J$ or $\omega_0^{(3,4)} = \omega_0^S - \pi J$ and $\omega_0^{(1,3)} = \omega_0^I + \pi J$ or $\omega_0^{(2,4)} = \omega_0^I - \pi J$ ($\omega_0^I = -\gamma_H B_0 < 0$ for ^1H , $\omega_0^S = -\gamma_N B_0 > 0$ for ^{15}N). The phase-cycling is: $\Phi_1 = y, y, -y, -y$, $\Phi_2 = y, -y, -y, y$, receiver $\Phi_{\text{rec}} = x, -x, x, -x$.

[0040] FIG. 5. Signal intensity of selected doublet components in boc-protected ^{15}N enriched glycine as a function of the spin-lock duration τ_2 (see Fig. 4). The experimental points correspond to cross-polarization from the high-frequency nitrogen transition $I_x^{(1,2)}$ to the low-frequency $I_x^{(1,3)}$ (■) or to the high-frequency $I_x^{(2,4)}$ (○) proton components (see schematic representation of the undecoupled HSQC multiplet and the corresponding resonance frequencies). Also displayed is a simulation of the (■) transfer taking into account relaxation within the three-dimensional subspace shown in Fig. 2, as described in the text. The experimental curves comprise 120 points, each of which was acquired with 16 scans at 303 K and 14 T.

[0041] FIG. 6. Fourier transform (ω_2) of the oscillatory τ_2 time dependence of ST-CP. Experimental points are calculated from the data of Fig. 5. Also displayed are the Fourier transforms of three simulations where the relaxation time T_2^{DOC} of the zero quantum coherence $I_y^{(2,3)}$ has been varied.

[0042] FIG. 7. Experimental T_2 decay curves of different nitrogen doublet components in human ubiquitin, obtained by inserting a Carr-Purcell-Meiboom-Gill (CPMG) multiple refocussing sequence before the ST-CP interval τ_2 of Fig. 4. For illustration two different amino acids (I36 and G76) are displayed. Filled symbols correspond to on-resonance cross-polarization from the high-frequency nitrogen transition $I_x^{(1,2)}$ to the low-frequency proton component $I_x^{(1,3)}$. Empty symbols correspond to off-resonance cross-polarization from the low-frequency nitrogen transition $I_x^{(3,4)}$ to the high-frequency proton component $I_x^{(2,4)}$. A schematic representation of the undecoupled HSQC multiplet indicates the on- and off-resonance components. Relaxation rates for both doublet components are thus taken from a single series of experiments with a total experiment time of 20 mins for each amide residue. The ST-CP cross-polarization transfer time was $\tau_2 = 26$ ms for field strengths of $\omega_1^I/2\pi = \omega_1^S/2\pi = 26$ Hz, which yields the maximum intensity for both the on-resonance and off-resonance components. Curve fits correspond to mono-exponential decays with T_2^S : (●) = 1.04 s, (○) = 457 ms, (■) = 306 ms, (D) = 87 ms. Each point was acquired with 64 scans at 303 K and 14 T.

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[0043]

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Claims

- 20 1. A method of nuclear magnetic resonance (=NMR) spectroscopy for achieving selective single-transition-to-single-transition cross-polarization (ST-CP) between two selected transitions in systems with two or more coupled spins by irradiating radio frequency (=RF) pulses to a sample in a homogeneous magnetic field in z-direction comprising the following steps:
- 25 1) Choose RF pulses suitably shaped for a selective excitation experiment with an RF amplitude that is weaker than the scalar coupling constant J_{IS} of the I and S spin systems;
- 2) Simultaneously stimulate a pair of connected transitions of the I and S spin systems;
- 30 3) Record the response for the said selective excitation experiment and observe different pairs of connected transitions.
2. The method according to claim 1, characterised in that the spin system is kept in isotropic or anisotropic solution during the experiment.
- 35 3. The method according to claim 1 or 2 characterised in that the intensities of the observed signals depend upon the chosen pair of connected transitions because of cross-correlated relaxation effects.
4. The method according to any one of the preceding claims, characterised by employing semi-selective cross-polarization during a preparation period τ_1 to transfer coherence through the scalar coupling from both doublet components of the selected source spin I to the two doublet components of the directly bound second spin S.
- 40 5. The method according to claim 4, characterised by using matched RF fields with amplitudes of the order of the heteronuclear coupling constant J_{IS} .
- 45 6. The method according to any one of the preceding claims, characterised by applying the following pulse sequence to the spin system I and S :
- a) irradiating a $\pi/2$ RF pulse acting on spins I with a phase-cycle $\Phi_1 = y, y, -y, -y$;
- 50 b) applying two RF fields with matched amplitudes of $\omega_1/2\pi = \omega_1^S/2\pi$ and carrier frequencies positioned at the midpoints of the I and the S doublets during a preparation time period $\tau_1 = J_{IS}^{-1}$;
- c) irradiating a $\pi/2$ RF pulse acting on spins S with a phase-cycle $\Phi_2 = y, -y, -y, y$
- d) destroying any residual transverse magnetization by applying a first spoil gradient pulse G_1 followed by a proton $\pi/2$ pulse and subsequently by applying a second spoil gradient pulse G_2 ;
- 55 e) irradiating a $\pi/2$ pulse acting on the spins S;
- f) applying two RF fields with matched amplitudes of $\omega_1/2\pi = \omega_1^S/2\pi$ for single-transition-to-single-transition cross-polarization, on-resonance with any chosen pair of connected transitions during a second time period τ_2 , the amplitudes of the RF fields being reduced with respect to the amplitudes in step b)

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[0038] FIG. 3. Simulated evolution of the density operator using the complete Hamiltonian of Eq. (3) and a relaxation superoperator as described in the text. RF field amplitudes are $\omega_1/2\pi = \omega_1'/2\pi = \omega_1^S/2\pi \approx 26$ Hz. The relaxation times used were $T_{1p}^I = 90$ ms, $T_{1p}^S = 415$ ms and $T_2^{DQC} = 30$ ms which are representative of an amino acid residue in the rigid core of a small protein. At $\tau = 24$ ms, one obtains a state with an optimum expectation value of

$$\langle I_x^{(2,4)} \rangle = 0.34.$$

[0039] FIG. 4. Pulse sequence used for coherence transfer by cross-polarization. The first semi-selective cross-polarization (CP) period has a duration $\tau_1 = J_{IS}^{-1} = 10.9$ ms and matched amplitudes of $\omega_1'/2\pi = \omega_1^S/2\pi \approx 70$ Hz. The carrier frequencies are positioned at the midpoints of the proton and nitrogen doublets. During the second, selective single-transition-to-single-transition cross-polarization (ST-CP) period τ_2 , the matched amplitudes are reduced to $\omega_1'/2\pi = \omega_1^S/2\pi \approx 26$ Hz and the carrier frequencies moved to any two connected transitions which appear in the undecoupled multiplet, e.g. $\omega_0^{(1,2)} = \omega_0^S + \pi J$ or $\omega_0^{(3,4)} = \omega_0^S - \pi J$ and $\omega_0^{(1,3)} = \omega_0^I + \pi J$ or $\omega_0^{(2,4)} = \omega_0^I - \pi J$ ($\omega_0^I = -\gamma_I B_0 < 0$ for 1H , $\omega_0^S = -\gamma_S B_0 > 0$ for ^{15}N). The phase-cycling is: $\Phi_1 = y, y, -y, -y, \Phi_2 = y, -y, -y, y$, receiver $\Phi_{rec} = x, -x, x, -x$.

[0040] FIG. 5. Signal intensity of selected doublet components in boc-protected ^{15}N enriched glycine as a function of the spin-lock duration τ_2 (see Fig. 4). The experimental points correspond to cross-polarization from the high-frequency nitrogen transition $I_x^{(1,2)}$ to the low-frequency $I_x^{(1,3)}$ (■) or to the high-frequency $I_x^{(2,4)}$ (○) proton components (see schematic representation of the undecoupled HSQC multiplet and the corresponding resonance frequencies). Also displayed is a simulation of the (■) transfer taking into account relaxation within the three-dimensional subspace shown in Fig. 2, as described in the text. The experimental curves comprise 120 points, each of which was acquired with 16 scans at 303 K and 14 T.

[0041] FIG. 6. Fourier transform (ω_2) of the oscillatory τ_2 time dependence of ST-CP. Experimental points are calculated from the data of Fig. 5. Also displayed are the Fourier transforms of three simulations where the relaxation time T_2^{DQC} of the zero quantum coherence $I_y^{(2,3)}$ has been varied.

[0042] FIG. 7. Experimental T_2 decay curves of different nitrogen doublet components in human ubiquitin, obtained by inserting a Carr-Purcell-Meiboom-Gill (CPMG) multiple refocussing sequence before the ST-CP interval τ_2 of Fig. 4. For illustration two different amino acids (I36 and G76) are displayed. Filled symbols correspond to on-resonance cross-polarization from the high-frequency nitrogen transition $I_x^{(1,2)}$ to the low-frequency proton component $I_x^{(1,3)}$. Empty symbols correspond to off-resonance cross-polarization from the low-frequency nitrogen transition $I_x^{(3,4)}$ to the high-frequency proton component $I_x^{(2,4)}$. A schematic representation of the undecoupled HSQC multiplet indicates the on- and off-resonance components. Relaxation rates for both doublet components are thus taken from a single series of experiments with a total experiment time of 20 mins for each amide residue. The ST-CP cross-polarization transfer time was $\tau_2 = 26$ ms for field strengths of $\omega_1'/2\pi = \omega_1^S/2\pi = 26$ Hz, which yields the maximum intensity for both the on-resonance and off-resonance components. Curve fits correspond to mono-exponential decays with T_2^S : (●) = 1.04 s, (○) = 457 ms, (■) = 306 ms, (D) = 87 ms. Each point was acquired with 64 scans at 303 K and 14 T.

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[0043]

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Claims

- 20 1. A method of nuclear magnetic resonance (=NMR) spectroscopy for achieving selective single-transition-to-single-transition cross-polarization (ST-CP) between two selected transitions in systems with two or more coupled spins by irradiating radio frequency (=RF) pulses to a sample in a homogeneous magnetic field in z-direction comprising the following steps:
- 25 1) Choose RF pulses suitably shaped for a selective excitation experiment with an RF amplitude that is weaker than the scalar coupling constant J_{IS} of the I and S spin systems;
- 2) Simultaneously stimulate a pair of connected transitions of the I and S spin systems;
- 30 3) Record the response for the said selective excitation experiment and observe different pairs of connected transitions.
2. The method according to claim 1, characterised in that the spin system is kept in isotropic or anisotropic solution during the experiment.
- 35 3. The method according to claim 1 or 2 characterised in that the intensities of the observed signals depend upon the chosen pair of connected transitions because of cross-correlated relaxation effects.
- 40 4. The method according to any one of the preceding claims, characterised by employing semi-selective cross-polarization during a preparation period τ_1 to transfer coherence through the scalar coupling from both doublet components of the selected source spin I to the two doublet components of the directly bound second spin S.
5. The method according to claim 4, characterised by using matched RF fields with amplitudes of the order of the heteronuclear coupling constant J_{IS} .
- 45 6. The method according to any one of the preceding claims, characterised by applying the following pulse sequence to the spin system I and S :
- a) irradiating a $\pi/2$ RF pulse acting on spins I with a phase-cycle $\Phi_1 = y, y, -y, -y$;
- 50 b) applying two RF fields with matched amplitudes of $\omega_1/2\pi = \omega_1^S/2\pi$ and carrier frequencies positioned at the midpoints of the I and the S doublets during a preparation time period $\tau_1 = J_{IS}^{-1}$;
- c) irradiating a $\pi/2$ RF pulse acting on spins S with a phase-cycle $\Phi_2 = y, -y, -y, y$;
- d) destroying any residual transverse magnetization by applying a first spoil gradient pulse G_1 followed by a proton $\pi/2$ pulse and subsequently by applying a second spoil gradient pulse G_2 ;
- 55 e) irradiating a $\pi/2$ pulse acting on the spins S;
- f) applying two RF fields with matched amplitudes of $\omega_1/2\pi = \omega_1^S/2\pi$ for single-transition-to-single-transition cross-polarization, on-resonance with any chosen pair of connected transitions during a second time period τ_2 , the amplitudes of the RF fields being reduced with respect to the amplitudes in step b)

g) using a receiver phase-cycling $\Phi_{\text{rec}} = x, -x, x -x$.

- 5 7. The method according to claim 6, characterised by inserting a suitable relaxation delay at a time point (a) before the ST-CP interval τ_2 to monitor the relaxation behaviour of the spin S coherences whilst observing the spin I doublet components.
- 10 8. The method according to claim 7, characterised by inserting a Carr-Purcell-Meiboom-Gill (CPMG) multiple refocussing sequence at time point (a) before the ST-CP interval τ_2 to monitor the relaxation behaviour of the spin S coherences whilst observing the spin I doublet components.
- 15 9. The method according to claim 8 using a Carr-Purcell-Meiboom-Gill (CPMG) spin echo sequence for T_2 measurements, characterised in that the intervals between the π pulses of the CPMG pulse train are set to $2\tau = n/J_{IS}$, where n is an integer.
- 20 10. The method according to claim 7, characterised by inserting a spin-lock field at time point (a) before the ST-CP interval τ_2 for $T_{1\rho}$ measurements.

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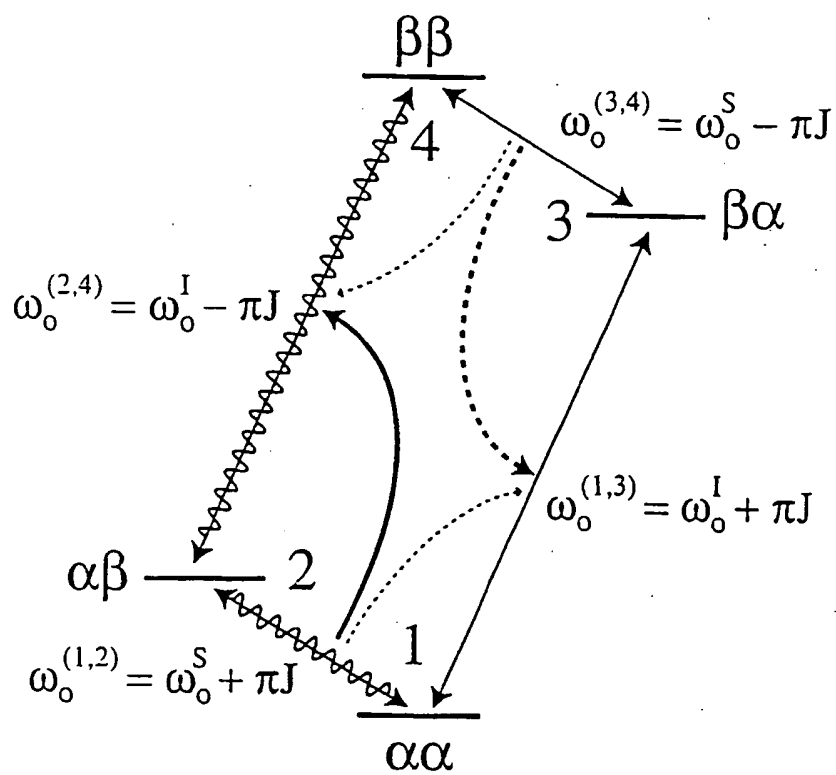


Figure 1

g) using a receiver phase-cycling $\Phi_{rec} = x, -x, x -x$.

7. The method according to claim 6, characterised by inserting a suitable relaxation delay at a time point (a) before the ST-CP interval τ_2 to monitor the relaxation behaviour of the spin S coherences whilst observing the spin I doublet components.
8. The method according to claim 7, characterised by inserting a Carr-Purcell-Meiboom-Gill (CPMG) multiple refocussing sequence at time point (a) before the ST-CP interval τ_2 to monitor the relaxation behaviour of the spin S coherences whilst observing the spin I doublet components.
9. The method according to claim 8 using a Carr-Purcell-Meiboom-Gill (CPMG) spin echo sequence for T_2 measurements, characterised in that the intervals between the π pulses of the CPMG pulse train are set to $2\tau = n/J_{IS}$, where n is an integer.
10. The method according to claim 7, characterised by inserting a spin-lock field at time point (a) before the ST-CP interval τ_2 for $T_{1\rho}$ measurements.

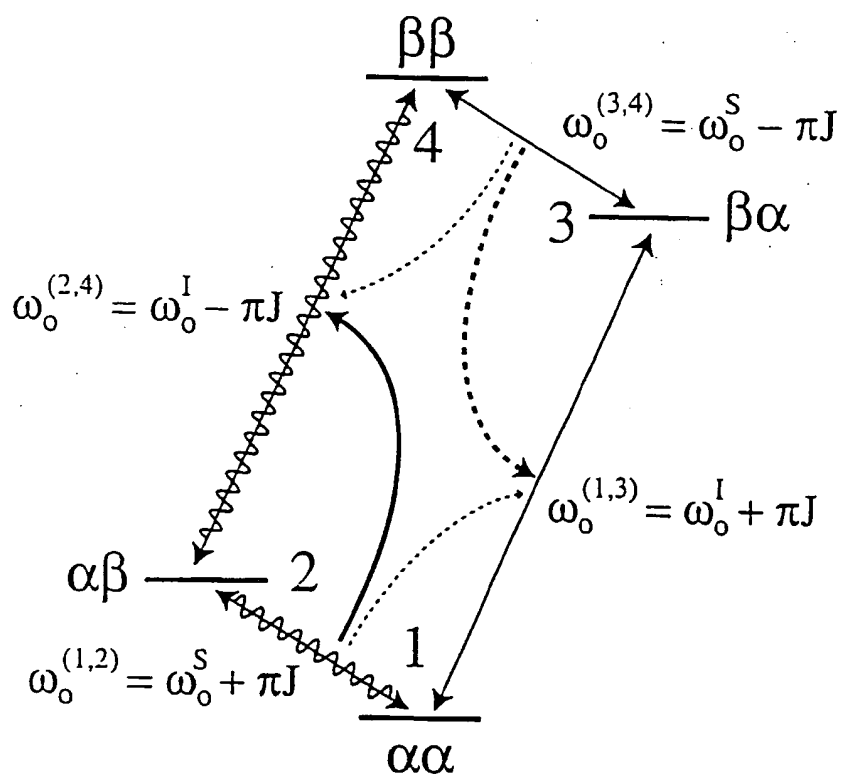


Figure 1

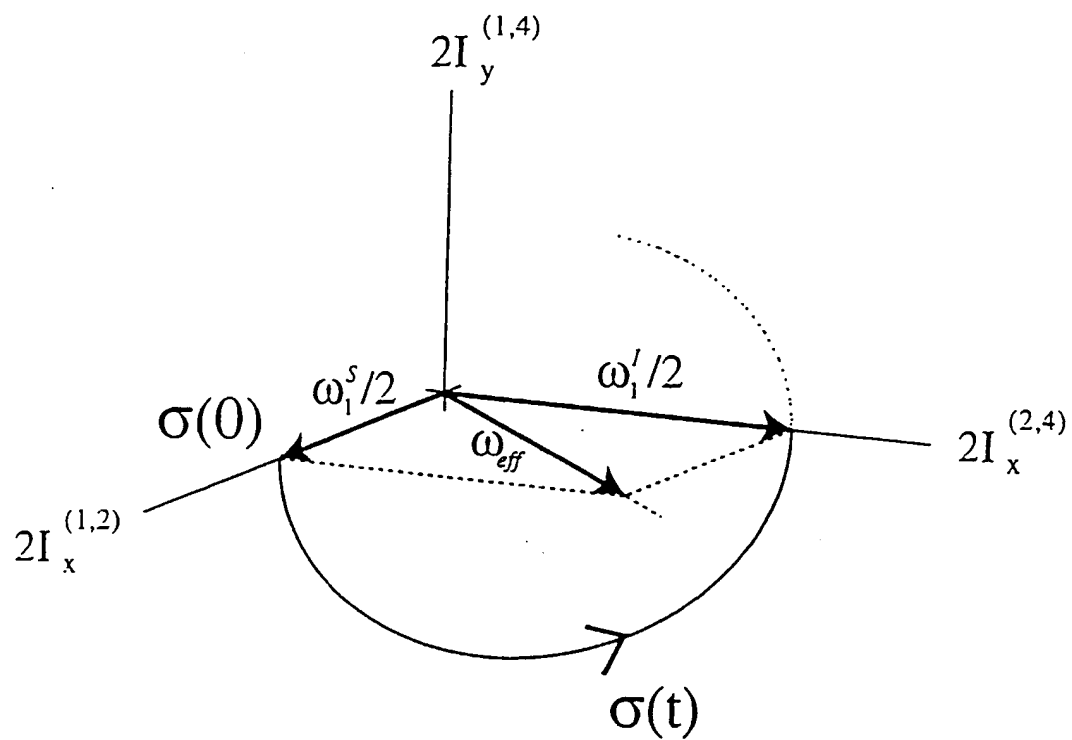


Figure 2

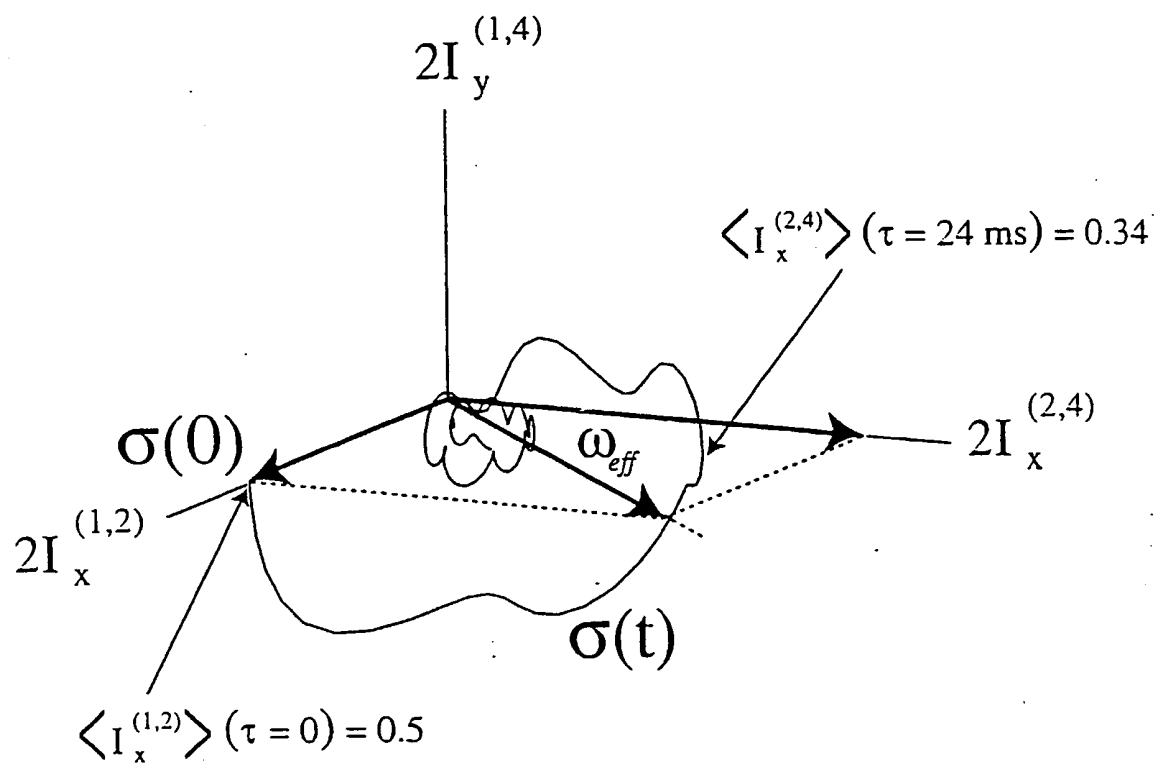


Figure 3

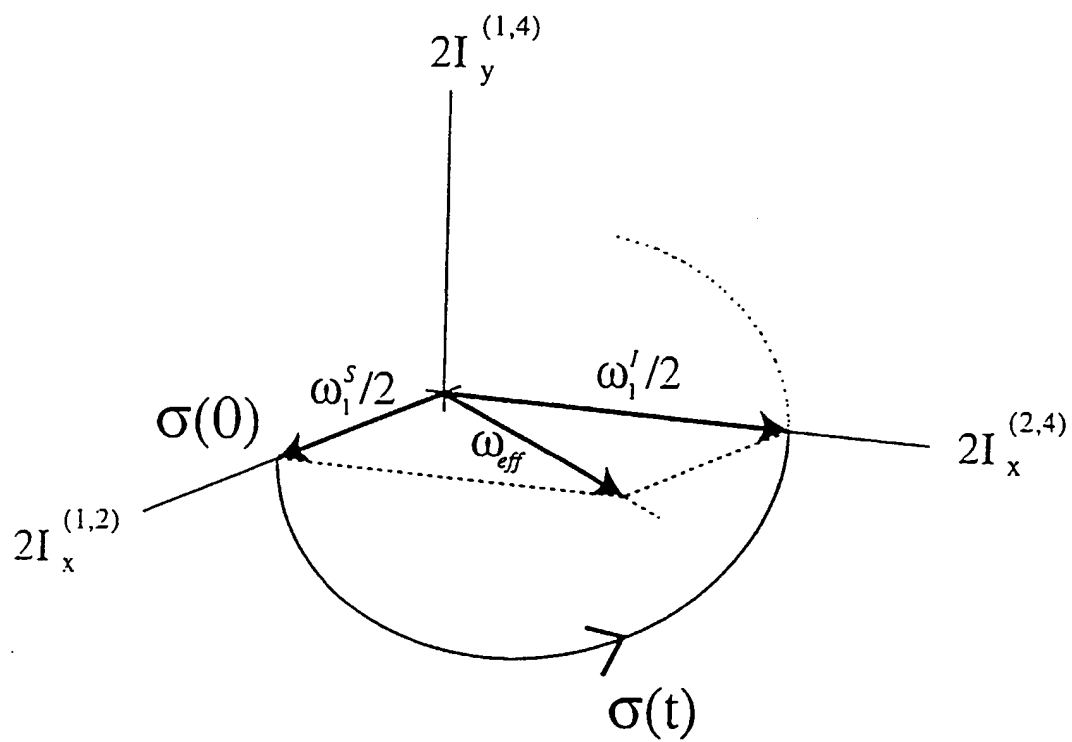
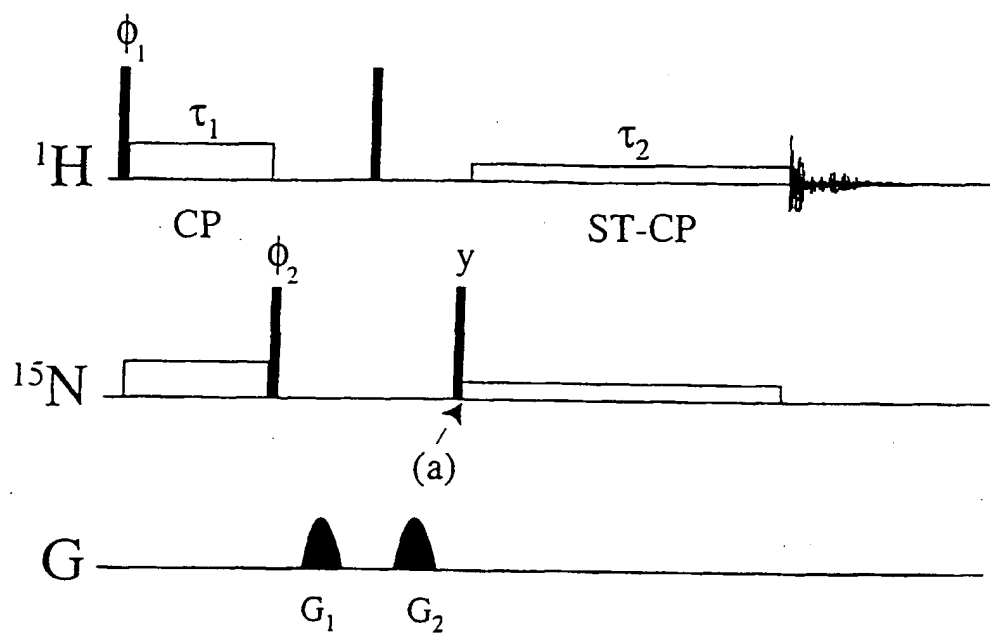
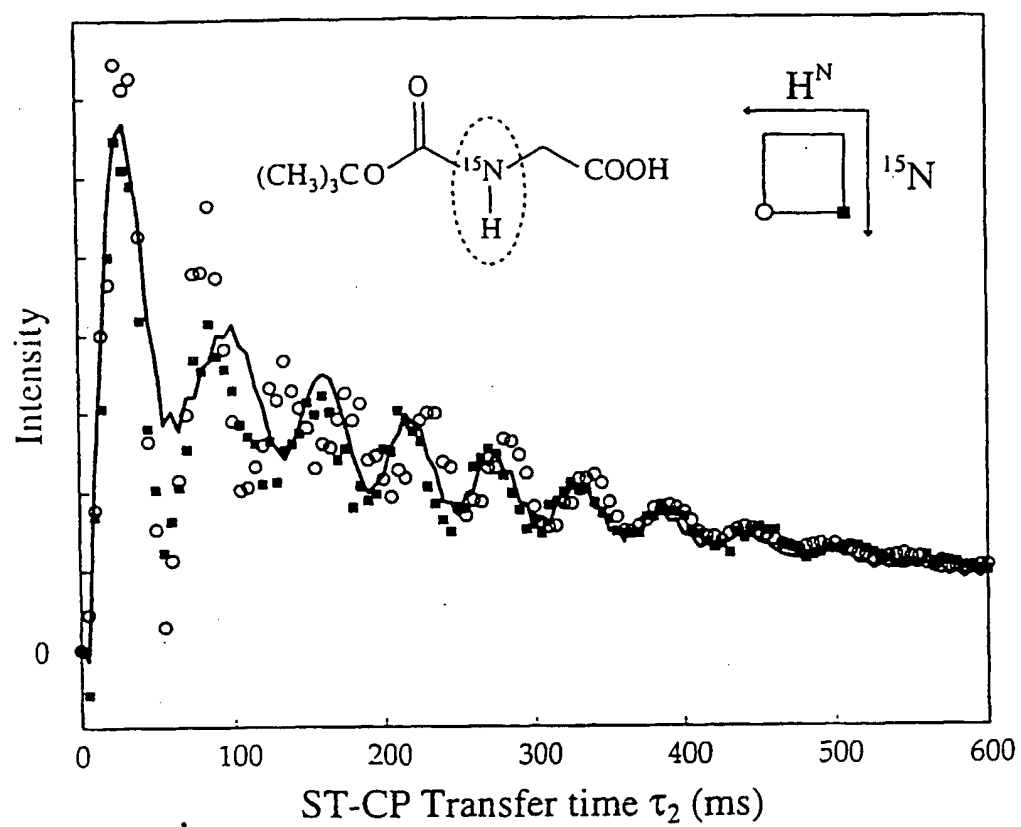
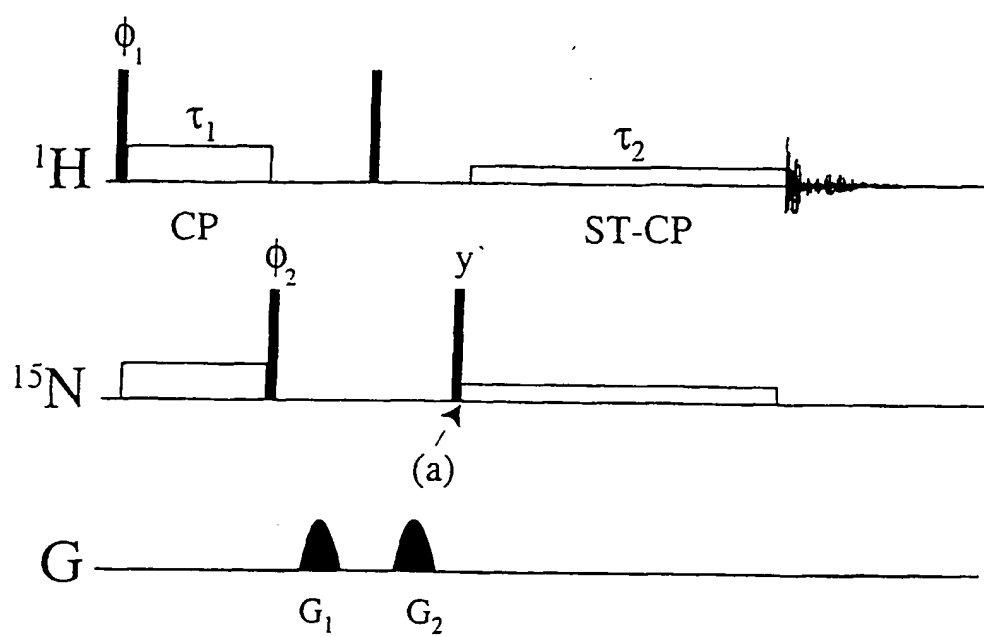
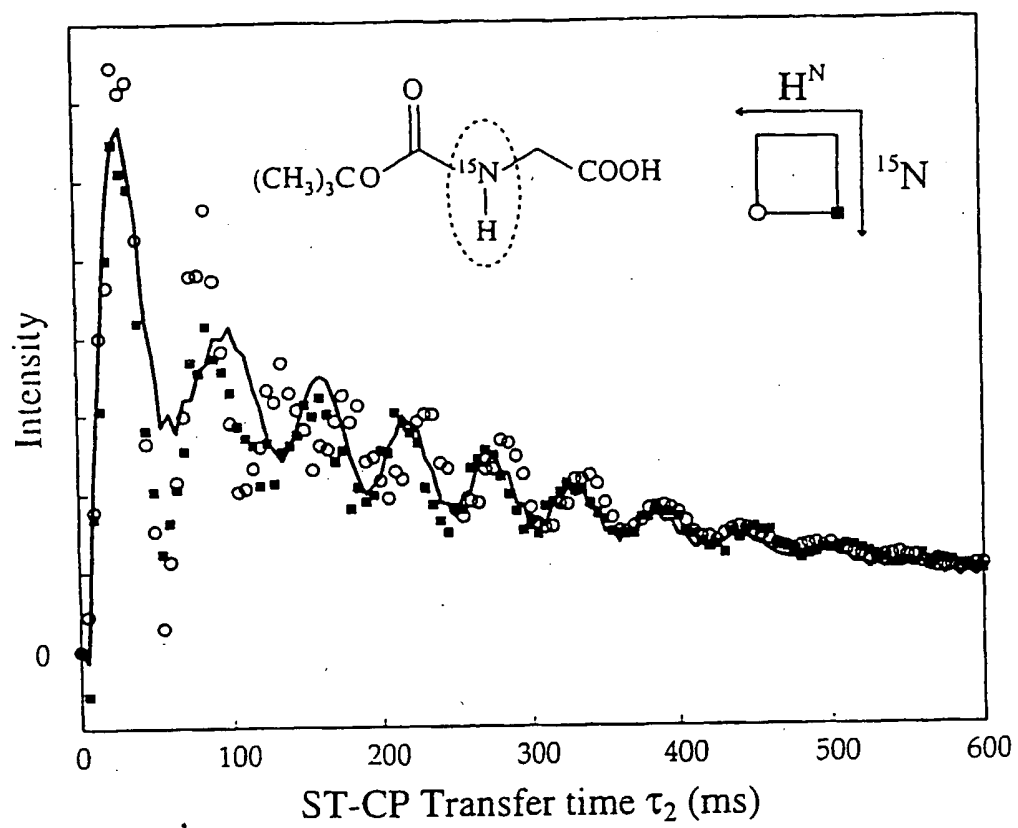


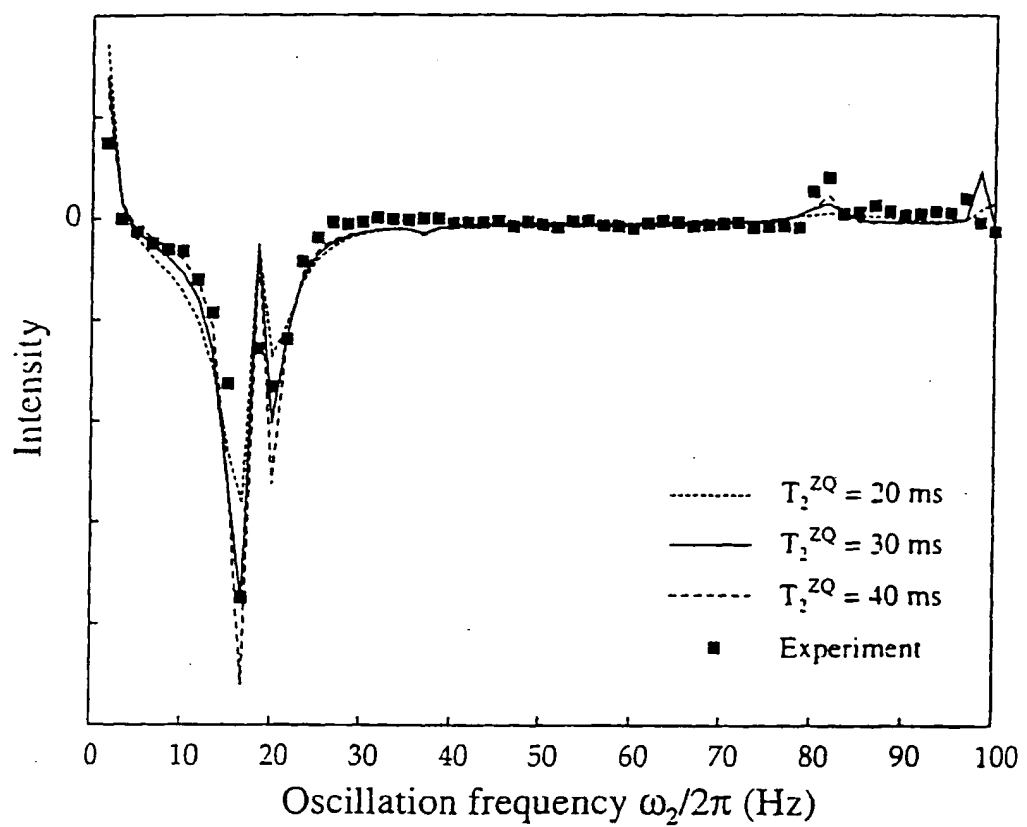
Figure 2

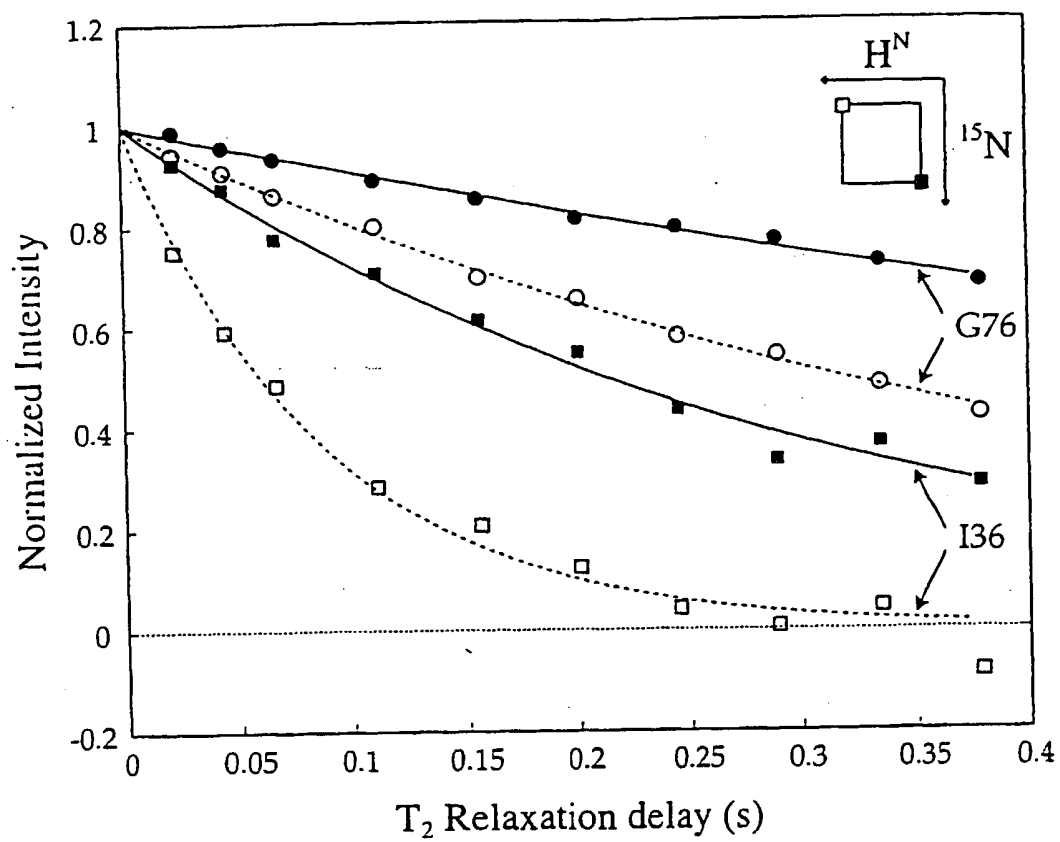
*Figure 4*

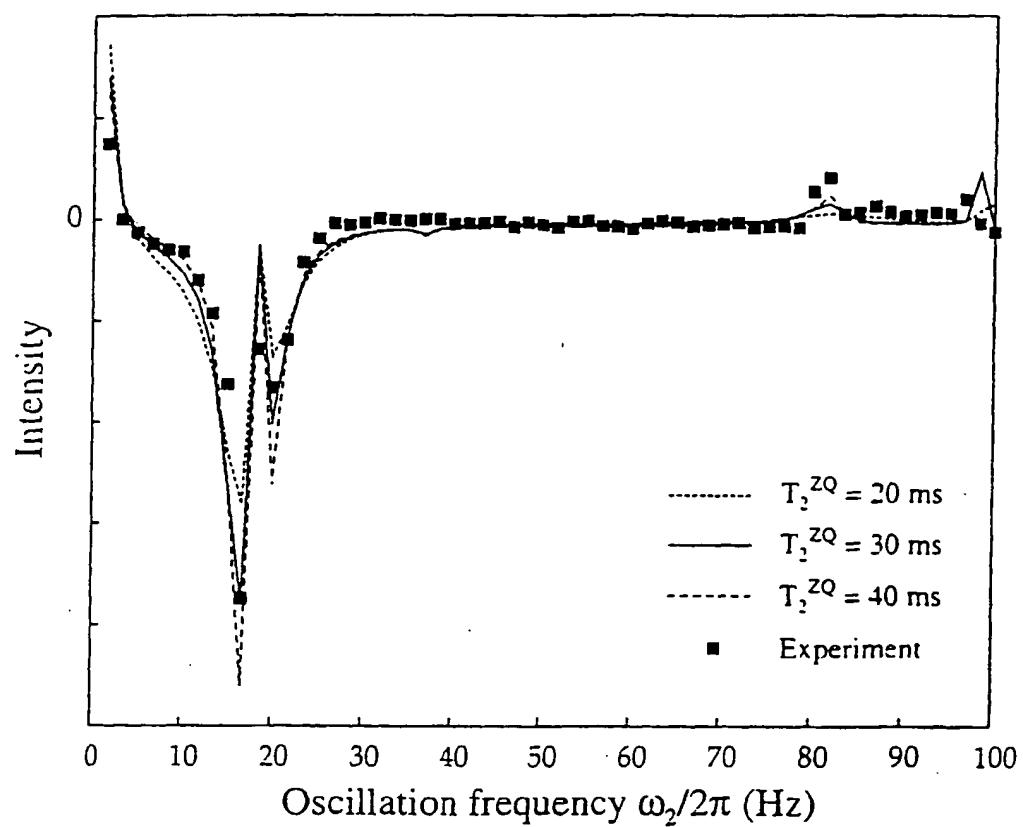
*Figure 5*

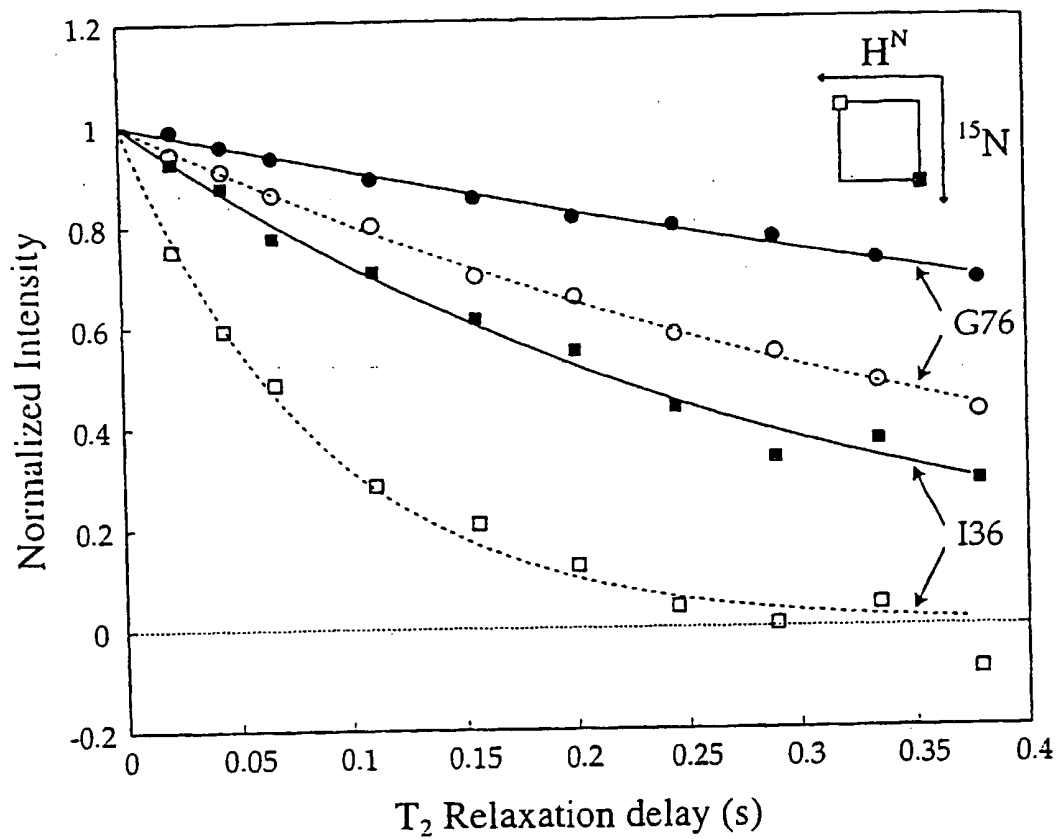
*Figure 4*

*Figure 5*

*Figure 6*

*Figure 7*

*Figure 6*

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EUROPEAN SEARCH REPORT

Application Number
EP 99 11 9845

DOCUMENTS CONSIDERED TO BE RELEVANT			
Category	Citation of document with indication, where appropriate, of relevant passages	Relevant to claim	CLASSIFICATION OF THE APPLICATION (Int.Cl.7)
A,D	ERMAKOV V L ET AL: "Coherence transfer in three-level systems: controlled violation of adiabaticity and antiparallel double resonant irradiation" JOURNAL OF CHEMICAL PHYSICS, 1 JULY 1995, USA, vol. 103, no. 1, pages 136-143, XP002129264 ISSN: 0021-9606	1-3,6	G01R33/46
A,D	CHIARPARIN E ET AL: "Selective cross-polarization in solution state NMR" MOLECULAR PHYSICS, 10 DEC. 1998, TAYLOR & FRANCIS, UK, vol. 95, no. 5, pages 759-767, XP000872079 ISSN: 0026-8976 * page 7 *	4-6	
			TECHNICAL FIELDS SEARCHED (Int.Cl.7)
			G01R
The present search report has been drawn up for all claims			
Place of search MUNICH		Date of completion of the search 4 February 2000	Examiner Lersch, W
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Application Number
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